A Simple Experiment to Explore Standing Waves in a Flexible Corrugated Sound Tube

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Sound tubes, pipes, and singing rods are used as musical instruments and as toys to perform amusing experiments. In particular, corrugated tubes present unique characteristics with respect to the sounds they can produce; that is why they have been studied so intensively, both at theoretical and experimental levels. Experimental studies usually involve expensive and sophisticated equipment that is out of reach of school laboratory facilities. In this paper we show how to investigate quantitatively the sounds produced by a flexible sound tube corrugated on the inside by using educational equipment readily available in school laboratories, such as the oscilloscope, the microphone, the anemometer, and the air pump. We show that it is possible for students to study the discontinuous spectrum of sounds produced by a flexible corrugated tube and go even further, computing the speed of sound in air with a simple experimental procedure.

Theory

When air flows across an open-ended tube, a standing air pressure wave is formed with a frequency that depends on the length of the tube. For an open-end air column, the fundamental standing wave (1st harmonic) has pressure nodes (therefore, air displacement antinodes) at both ends and a pressure antinode (displacement node) in the middle of the tube. The length \( L \) of the tube corresponds, in this case, to \( L = \frac{\lambda_0}{2} \), where \( \lambda_0 \) is the wavelength of the standing wave. The corresponding frequency \( f_0 \) of the standing wave can be measured with an oscilloscope, capturing the sound with a microphone. High-order standing waves can also be formed, each one with a frequency that is a multiple of the fundamental wave, as shown in Fig. 1.

Relating \( \lambda_n \) with \( f_n \), where \( n \) is the order of the harmonic standing wave, we can calculate the speed of sound in air,

\[
\nu_{\text{sound}} = 2f_0. \tag{1}
\]

Experimental procedure and results

Sound tubes (and pipes) produce sounds when air goes through them. In flexible sound tubes corrugated on the inside, as the air flows over one ridge and then over another, it tumbles into a vortex. When the frequency of the vortex matches one of the natural resonant frequencies of the tube, sound is amplified. The faster the air flows through the tube, the higher the pitch of the sound produced by the vortex.

When a flexible sound tube is spun in circular motion, there is a decrease of the air pressure \( P \) at the free extremity,
due to the speed it reaches, according to Bernoulli’s law,

$$P + \rho gh + \frac{1}{2} \rho v^2 = \text{constant},$$  \hspace{1cm} (2)

where $P$ is the fluid pressure (or energy per unit volume), $ho gh$ is the potential energy per unit volume, and $\frac{1}{2} \rho v^2$ is the kinetic energy per unit volume. Therefore, a differential in air pressure is created between the extremity of the tube held by the hand (higher air pressure) and the free moving extremity (lower air pressure) (see Fig. 2), pushing air through the tube: the faster the tube spins, the higher the speed of air across it.

Such a situation was simulated in the laboratory by pumping air through one extremity of the tube so we could control the flux of air crossing it; at the other extremity, a digital anemometer Kestrel 1000 measured the speed of air, and a standard microphone captured the sound wave produced inside the tube. The electrical signal was then visualized in a digital oscilloscope Nicolet 310 and the period of the sound wave was measured, allowing us to know the corresponding frequency of the signal (see Fig. 3). Details of the experimental setup are shown in Fig. 4.

Results in Fig. 3 could also be obtained by replacing the oscilloscope in our setup in Fig. 4 with data acquisition software such as MacScope, PASCO Data Studio, Vernier Logger Pro, or Pico Technology’s PicoScope, among others.

Figure 5 shows the frequency of the sound waves produced as a function of the air speed $v_{\text{air}}$ across the tube. It is easy to see that the frequency does not change continuously with $v_{\text{air}}$, but instead a well-defined sequence of steps is detected, corresponding to the different sounds that were heard.

By plotting the average frequency of each step as a function of the number of frequency steps detected ($N$), we were able to determine the fundamental frequency $f_0$ corresponding to the first harmonic. Results are resumed in Fig. 6.

A careful analysis of the linear fit to the experimental data in Fig. 6 shows that the first step in Fig. 5 corresponds, in fact, to the second harmonic of the standing waves in the sound tube: the first harmonic was actually calculated from the slope of this straight line and is $f_0 = (214 \pm 3)$ Hz. This sound could not be detected during the experiment; a similar situation was also reported in Ref. 4.

For an ideal open-ended tube, the speed of sound in air can be calculated using Eq. (1). However, as the air is moving in and out of the tube, the pressure variations do not drop to zero right at the open ends of the tube, but rather a small distance beyond. Therefore, the effective acoustic length $L'$ is slightly greater than the physical length $L$ of the tube. For a cylindrical tube of radius $r$, a correction for the additional length must be introduced:4,14

$$L' = L + 1.22r.$$  \hspace{1cm} (3)

As the physical length of the sound tube is $L = 75.0$ cm and its radius $r = 1.5$ cm, the acoustic length according to Eq. (3) is
The presence of a discontinuous spectrum of frequencies, characterized by well-defined frequency steps with an increasing sound pitch as the tube spins, is a clear indication of the formation of standing waves inside the sound tube. The analysis of experimental data supports motivation for a discussion about harmonics in sound waves and allows the calculation of the speed of sound in air with a reasonable accuracy.

References